

# Complex Economic Systems in Macro & Finance

## Lecture I: Complex Systems, Nonlinear Dynamics, Behavioral Rationality and Heterogeneous Expectations

Cars Hommes

CeNDEF, University of Amsterdam

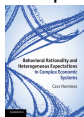
CEF 2015 Workshop  
Taipei, Taiwan, June 2015

# Outline of the Workshop: 4 Lectures

- 1 Complex systems, nonlinear dynamics, behavioral rationality, heterogeneous expectations
- 2 Asset pricing model with heterogeneous beliefs
- 3 Empirical validation and Laboratory Experiments
- 4 Behavioral macroeconomics with heterogeneous expectations

# Some References

- Hommes, C.H., (2013), Behavioral Rationality and Heterogeneous Expectations in Complex Economic Systems, Cambridge.



- complex nonlinear dynamics: Chapters 2 (1-D) and 3 (2-D + higher D)
- nonlinear cobweb model with homogeneous expectations (Chapter 4)
- cobweb model with heterogeneous expectations (rational versus naive) (Chapter 5)
- asset pricing model with heterogeneous expectations (Chapter 6)
- empirical validation (Chapter 7)
- laboratory experiments (Chapter 8)
- some more recent papers on behavioral macro and housing market

# What is a Complex System

## Characteristic Features

- **nonlinear**, complicated dynamics (multiple steady states, cycles, chaos)
- **interacting** particles/agents
- **emergent macro behavior** may be different from **individual micro behavior**
- **heterogeneous** agents, different from representative agent
- **complex socio-economic systems**: the particles can think need a theory of adaptive behavior
- simple behavioral **heuristics** instead of perfect rationality

# Outline

- 1 The 1-D quadratic map
- 2 Bifurcations & Chaos
- 3 Strange Attractors
- 4 Hopf bifurcation
- 5 Homoclinic orbits
- 6 Implications Nonlinear Dynamics for Economics

# Quadratic Map

Example of one-dimensional system:

$$x_{t+1} = f_{\lambda}(x_t) = \lambda x_t(1 - x_t) \quad (2.2)$$

- ① initial state  $x_0 \in [0, 1]$
- ② **parameter**  $\lambda$ ,  $0 \leq \lambda \leq 4$ .
- ③ **Problem:** what do the **orbits** look like?

# Convergence to a steady state

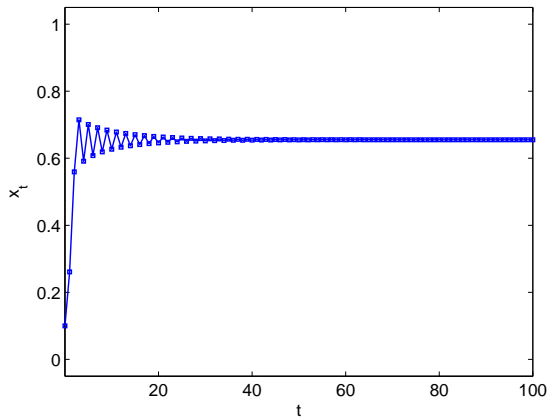


Figure :  $\lambda = 2.9$  and  $x_0 = 0.1$

# Convergence to a 2-cycle

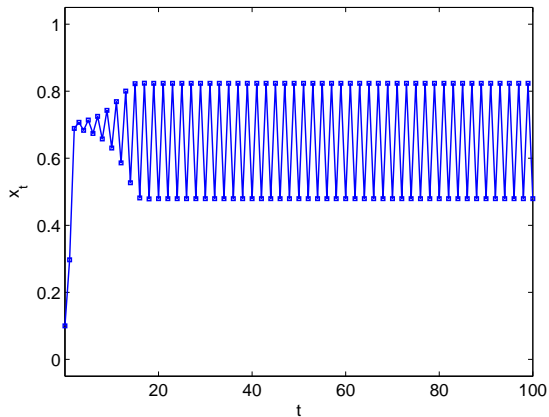


Figure :  $\lambda = 3.3$  and  $x_0 = 0.1$



# Convergence to a 4-cycle

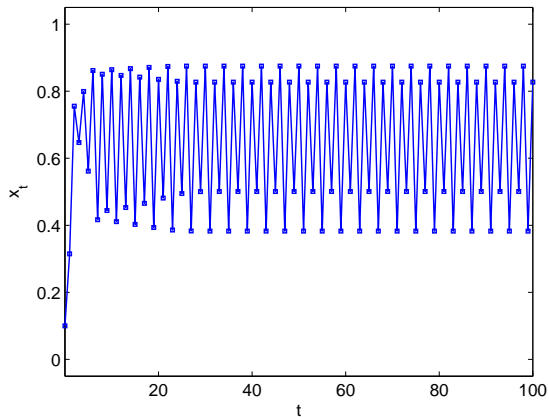


Figure :  $\lambda = 3.5$  and  $x_0 = 0.1$

# Convergence to a 3-cycle

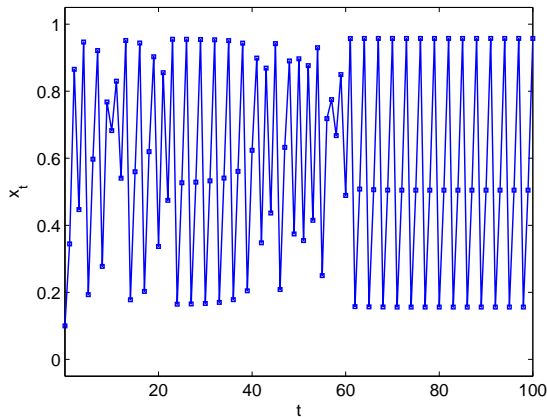


Figure :  $\lambda = 3.83$  and  $x_0 = 0.1$

# Sensitive dependence on initial conditions

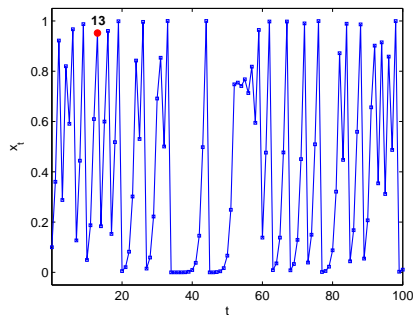
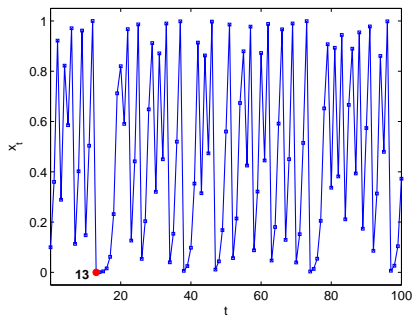


Figure :  $\lambda = 4$  and  $x_0 = 0.1$  (left) and (f)  $\lambda = 4$  and  $x_0 = 0.1001$  (right).

# Periodic Orbits and Stability

A point  $x$  is called a **periodic point with period  $k$**  if

$$f^k(x) = x \quad \text{and} \quad f^i(x) \neq x, \quad 0 < i < k.$$

(Note: periodic point with period  $k$  is fixed point of  $k$ -th iterate  $f^k$ )

$\{x_1, x_2, \dots, x_k\} = \{x_1, f(x_1), f^2(x_1), \dots, f^{k-1}(x_1)\}$  **periodic orbit** or  **$k$ -cycle**.

If  $x_i$  stable fixed point of  $f^k$ , then  $\{x_1, x_2, \dots, x_k\}$  **stable periodic orbit**;  
from the chain rule we have

$$\begin{aligned} (f^k)'(x_i) = (f^k)'(x_1) &= f'(f^{k-1}(x_1)) \cdot f'(f^{k-2}(x_1)) \dots f'(f(x_1)) \cdot f'(x_1) \\ &= \prod_{i=0}^{k-1} f'(f^i(x_1)). \end{aligned}$$

(Note:  $(f^k)'(x_j)$  is the product of derivatives along the orbit)

# Aperiodic Point

A point  $x$  is called an **aperiodic point** if

- 1 the orbit of  $x$  is bounded,
- 2 the orbit of  $x$  is not periodic, and
- 3 the orbit of  $x$  does not converge to a periodic orbit

# Bifurcation diagram of the quadratic map

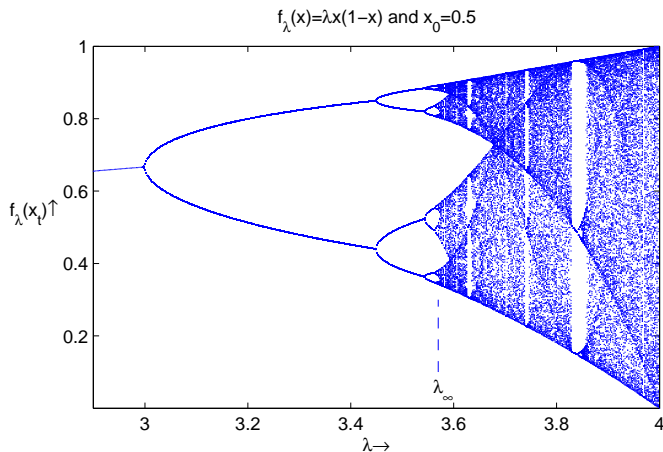
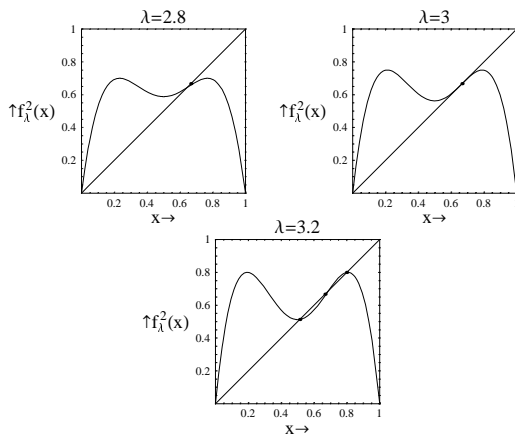


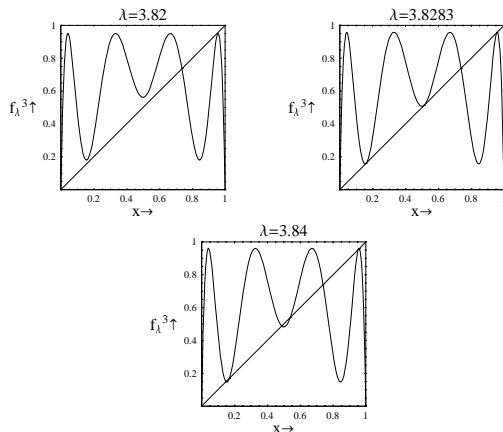
Figure : Bifurcation diagram of the quadratic map.

# Period doubling bifurcation at $\lambda = 3$



**Figure :** Graphs of the second iterate  $f^2$  for three different  $\lambda$ -values close to the period-doubling bifurcation at  $\lambda = 3$ .

# Creation of a 3-cycle by tangent bifurcation



**Figure :** Creation of 3-cycle by tangent bifurcation at  $\lambda \approx 3.8283$ .



# Definition of topological chaos

The dynamics of a difference equation  $x_{t+1} = f(x_t)$  is called **(topologically) chaotic** if the following three properties are satisfied:

- ① There exists an infinite set  $P$  of (unstable) periodic points with different periods.
- ② There exists an uncountable set  $S$  of aperiodic points (i.e. pointset whose orbits are bounded, not periodic and not converging to a periodic orbit).
- ③  $f$  has sensitive dependence on initial conditions w.r.t.  $\Lambda = P \cup S$ , that is, there exists a positive distance  $C$  such that for all initial states  $x_0 \in \Lambda$  and any  $\varepsilon$ -neighbourhood  $U$  of  $x_0$ , there exists an initial state  $y_0 \in \Lambda \cap U$  and a time  $T > 0$  such that the distance  $d(x_T, y_T) = d(f^T(x_0), f^T(y_0)) > C$ .

# Example of Chaos: quadratic map for $\lambda = 4$

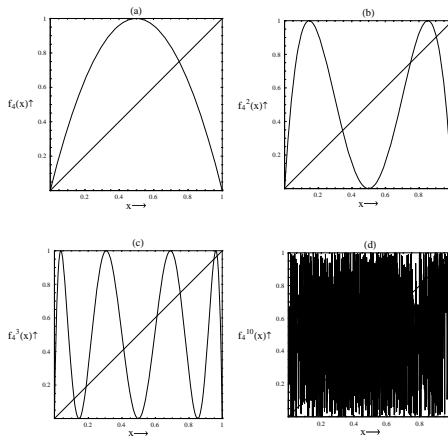


Figure : Graphs of (a)  $f_4(x) = 4x(1-x)$ , (b)  $f_4^2$ , (c)  $f_4^3$  and (d)  $f_4^{10}$ .

# Properties of quadratic map $f_4$

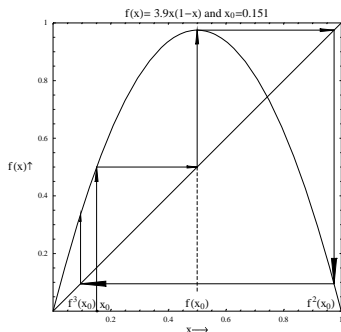
It can be shown that for any  $n$ , the graph of  $f_4^n$  has the following properties:

- ①  $f_4^n$  has  $2^{n-1}$  maxima equal to 1 and  $2^{n-1} + 1$  minima equal to 0 (including minima at  $x = 0$  and  $x = 1$ ).
- ②  $f_4^n$  'oscillates'  $2^{n-1}$  times on the interval  $[0, 1]$ .
- ③ the map  $f_4^n$  has  $2^n$  fixed points.
- ④ for any interval  $I$  of arbitrarily small length  $\varepsilon$ , there exists an  $N > 0$  such that  $I$  contains points  $x, y$  with  $f_4^N(x) = 0$  and  $f_4^N(y) = 1$ .

# Period three implies chaos

## Theorem 1

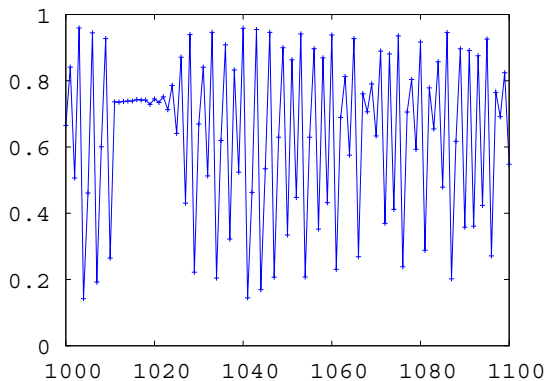
("Period 3 implies Chaos", Li & Yorke [1975]). Let  $x_{t+1} = f(x_t)$  be a 1-D difference equation with  $f$  a continuous map. If there exist a point  $x_0$  such that  $f^3(x_0) \leq x_0 < f(x_0) < f^2(x_0)$  (or with  $>$  instead of  $<$ ) then the dynamics is topologically chaotic



# Topological Chaos with Noise

Quadratic map with small noise

$$x_{t+1} = 3.83x_t(1 - x_t)$$



# Lyapunov exponents

- The **Lyapunov exponent** is defined as  

$$\lambda(x_0) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=0}^{n-1} \ln(|f'(f^i(x_0))|).$$
- Derivation:

$$|f^n(x_0 + \delta) - f^n(x_0)| \approx |(f^n)'(x_0)\delta| = e^{n\lambda(x_0)} |\delta|$$

$$\Leftrightarrow e^{n\lambda(x_0)} = |(f^n)'(x_0)| \Rightarrow \lambda(x_0) = \frac{1}{n} \ln(|(f^n)'(x_0)|).$$

- The Lyapunov exponent measures the **average rate of divergence of nearby initial states**. It is the average of the logs of the absolute values of the derivative along the orbit.

# Lyapunov exponent plot of the quadratic map

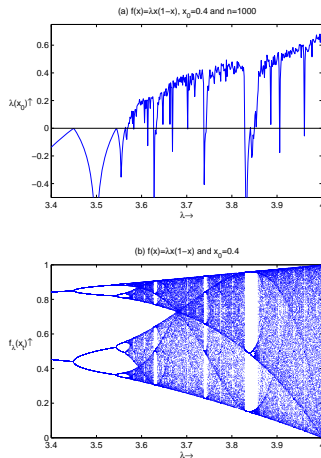


Figure : Lyapunov exponent  $L$  as a function of the parameter  $\lambda$ .

## Two-dimensional (2-D) systems

$$(x_{t+1}, y_{t+1}) = F_{\lambda}(x_t, y_t),$$

$F_{\lambda}$  nonlinear 2-D map and  $\lambda$  is a parameter.

The **orbit** with **initial state**  $(x_0, y_0)$  is the set

$$\{(x_0, y_0), (x_1, y_1), (x_2, y_2), \dots\} = \{(x_0, y_0), F_{\lambda}(x_0, y_0), F_{\lambda}^2(x_0, y_0), \dots\}.$$

**Problem:** what do these orbits look like and how does it depend on initial states and parameters?

**Example:** Hénon map:

$$\begin{aligned} x_{t+1} &= 1 - ax_t^2 + y_t \\ y_{t+1} &= bx_t, \end{aligned}$$

where  $a$  and  $b$  are parameters.

(special case  $b = 0$  yields 1-D quadratic map)



# Strange attractor of the Hénon map

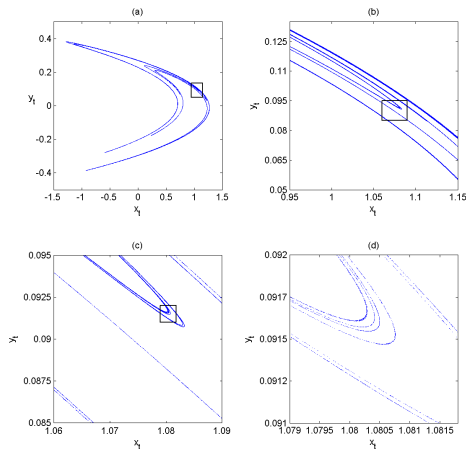
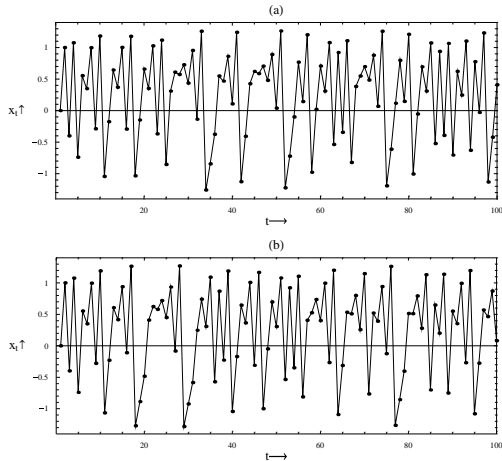


Figure : The strange attractor for the Hénon map  $H_{a,b}$  with  $a = 1.4$  and  $b = 0.3$ .

# Chaotic time-series and SDIC



**Figure :** Chaotic time series and sensitive dependence for the Hénon map  $H_{a,b}$  with  $a = 1.4$  and  $b = 0.3$ . (a)  $(x_0, y_0) = (0, 0)$  and (b)  $(x_0, y_0) = (0.001, 0)$ .

# Attractor and Strange Attractor

An **attractor** of a  $K$ -dimensional system  $X_{t+1} = F(X_t)$  is a compact set  $A$  with the following properties:

- ① The set  $A$  is **invariant**, i.e.  $F(A) \subset A$ .
- ② The set  $A$  is **invariant**, i.e. there exists an open neighborhood  $U$  of  $A$  (i.e.  $A \subset U$ ), such that all initial states  $X_0$  converge to the attractor  $A$ , i.e. for all  $X_0 \in U$ ,  $\lim_{n \rightarrow \infty} \text{dist}(F^n(X_0), A) = 0$ .
- ③ There exists an initial state  $X_0 \in A$  for which the orbit is **dense** in  $A$ .

An attractor  $A$  is called a **strange attractor** of the  $N$ -dimensional dynamical system  $x_{t+1} = F(x_t)$ , if the map  $F$  has sensitive dependence w.r.t. the set of initial states converging to  $A$ .

# The delayed logistic map

- Delayed logistic map:  $N_{t+1} = aN_t(1 - N_{t-1})$ .
- Equivalently ( $x_t = N_t$  and  $y_t = N_{t-1}$ ):

$$\begin{aligned}x_{t+1} &= y_t \\ y_{t+1} &= ay_t(1 - x_t).\end{aligned}$$

- **steady states**

$$(x_1, y_1) = (0, 0) \quad \text{and} \quad (x_2, y_2) = \left(\frac{a-1}{a}, \frac{a-1}{a}\right).$$

- The **eigenvalues** of the system are  $\lambda_1 = \frac{1}{2} - \frac{1}{2}\sqrt{5-4a}$  and  $\lambda_2 = \frac{1}{2} + \frac{1}{2}\sqrt{5-4a}$ .

# Dynamical properties of the delayed logistic map

The eigenvalues  $\lambda_1$  and  $\lambda_2$  of  $JF_a(\frac{a-1}{a}, \frac{a-1}{a})$  satisfy the following properties:

- $0 \leq a < 1$ : real eigenvalues with  $-1 < \lambda_1 < 1 < \lambda_2$ , so  $(\frac{a-1}{a}, \frac{a-1}{a})$  is a **saddle**.
- $1 < a < \frac{5}{4}$ : real eigenvalues with  $0 < \lambda_1 < \lambda_2 < 1$ , so  $(\frac{a-1}{a}, \frac{a-1}{a})$  is attracting (**stable node**).
- $\frac{5}{4} < a < 2$ : complex eigenvalues with  $\lambda_1 \lambda_2 = a - 1 < 1$ , so  $(\frac{a-1}{a}, \frac{a-1}{a})$  is a **stable focus**.
- $a > 2$ : complex eigenvalues with  $\lambda_1 \lambda_2 = a - 1 > 1$ , so  $(\frac{a-1}{a}, \frac{a-1}{a})$  is an **unstable focus**.
- **Hopf bifurcation** (or Neimark-Sacker) for  $a = 2$   
complex eigenvalues on the unit circle

# Attractors delayed logistic map

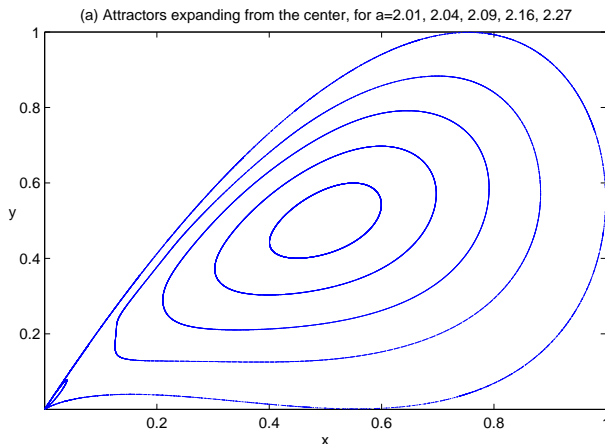


Figure : (a) Attractors for the logistic delayed equation for different  $a$ -values after the Hopf bifurcation.

# Strange attractor of the delayed logistic map

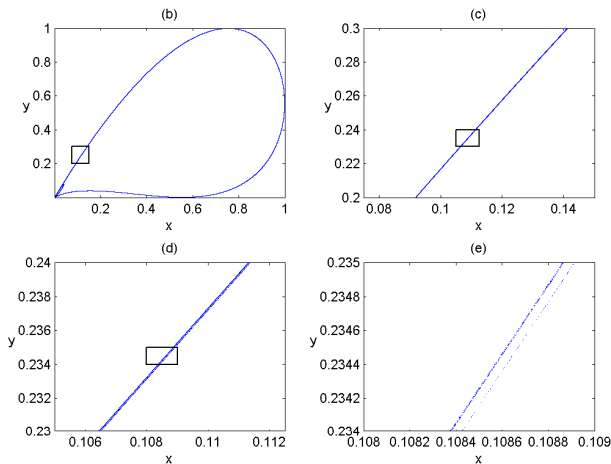


Figure : (b-e) The strange attractor for  $a = 2.27$  and some enlargements.

# Time series for the logistic delayed equation

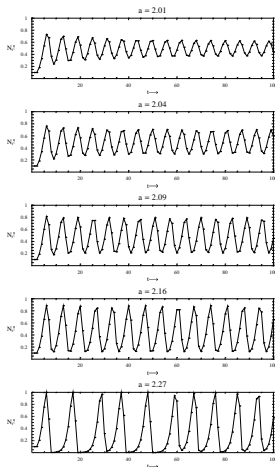
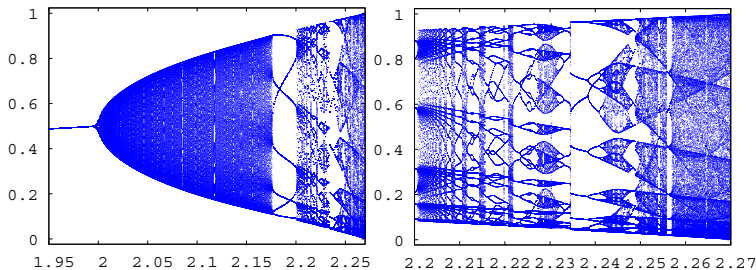


Figure : Time series for for different values of the parameter  $a$ .



# Bifurcation diagrams for logistic delayed equation



**Figure :** Hopf-bifurcation and breaking of an invariant circle bifurcation route to chaos

# Types of Local Bifurcations

- 1  $\lambda = +1$ :
  - **saddle-node bifurcation**: two new steady states, a saddle and a node;
  - **pitchfork bifurcation**: one steady state becomes unstable and two new stable steady states;
  - **transcritical bifurcation**; two steady states collide and exchange stability;
- 2  $\lambda = -1$ :
  - **period-doubling bifurcation**: steady state loses stability and new stable 2-cycle;
- 3 a pair of complex eigenvalues  $\lambda_1$  and  $\lambda_2$  on the unit circle, i.e.  $|\lambda_1 \lambda_2| = 1$  :
  - **Hopf bifurcation**: steady state becomes an unstable focus and an attracting invariant circle emerges with (quasi-)periodic dynamics

# Local (un)stable manifolds

Let  $p$  be a fixed point of the 2-D map  $F$ .

The **local stable manifold** and **local unstable manifold** of  $p$  are defined as

$$W_{loc}^s(p) = \{x \in U \mid \lim_{n \rightarrow \infty} F^n(x) = p\} \quad (1)$$

$$W_{loc}^u(p) = \{x \in U \mid \lim_{n \rightarrow -\infty} F^n(x) = p\} \quad (2)$$

where  $U$  is some small neighbourhood of  $p$ .

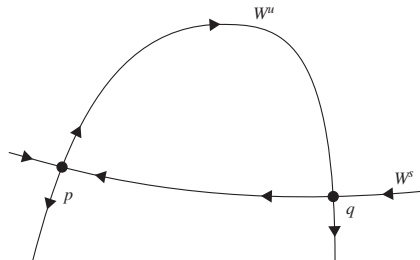
# Global(un)stable manifolds

The **global stable manifold** and the **global unstable manifold** are now defined as

$$W^s(p) = \bigcup_{n=0}^{\infty} F^{-n}(W_{loc}^s) \quad (3)$$

$$W^u(p) = \bigcup_{n=0}^{\infty} F^n(W_{loc}^u). \quad (4)$$

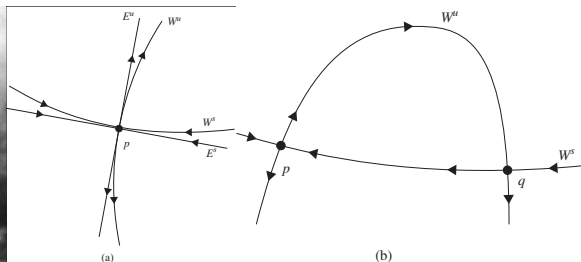
# Homoclinic point



(b)

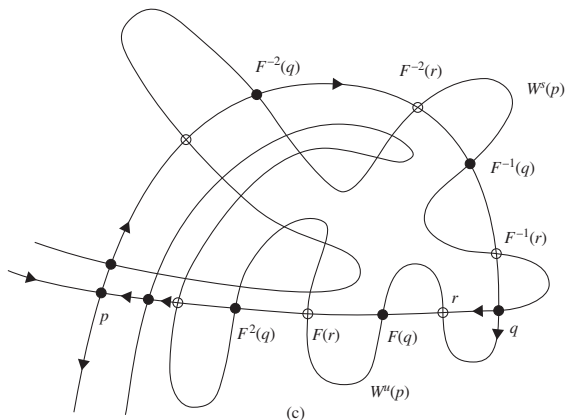
A point  $q$  is called a **homoclinic point** if  $q \neq p$  and  $q$  is an intersection point of the stable and unstable manifolds of the saddle point  $p$ , that is,  $q \in W^s(p) \cap W^u(p)$ .

Henry Poincaré, ca. 1890: motion in the **three body problem** is unpredictable and chaotic



2-D map of suitable plane section has **homoclinic orbit**

# Henry Poincaré, ca. 1890: homoclinic orbit implies sensitive dependence on initial conditions



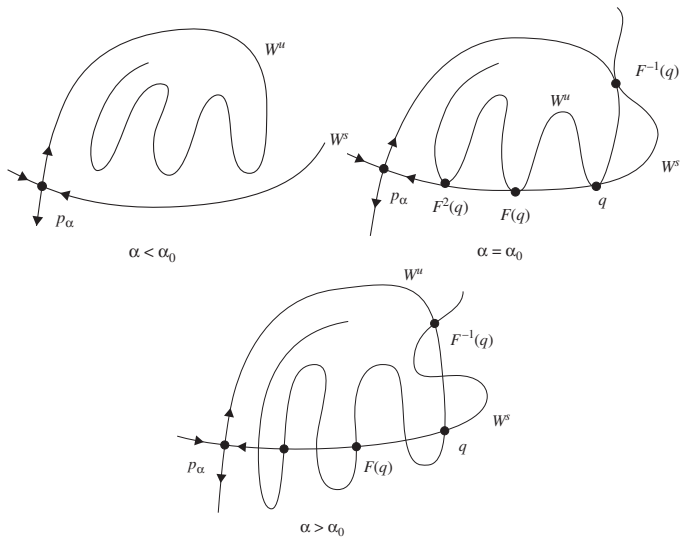
# Homoclinic bifurcation

We say that  $F_\alpha$  has a **homoclinic bifurcation**, associated to the saddle point  $p_\alpha$ , at  $\alpha = \alpha_0$ , if

- 1 for  $\alpha < \alpha_0$ ,  $W^s(p_\alpha)$  and  $W^u(p_\alpha)$  have no intersection point  $q \neq p$ ;
- 2 for  $\alpha = \alpha_0$ ,  $W^s(p_\alpha)$  and  $W^u(p_\alpha)$  have a point of homoclinic tangency;
- 3 for  $\alpha > \alpha_0$ ,  $W^s(p_\alpha)$  and  $W^u(p_\alpha)$  have a transversal homoclinic intersection point.



# Homoclinic bifurcation



# Implications Nonlinear Dynamics for Economics

- The fact that (simple) nonlinear systems exhibit **complex dynamics** calls for **reservations about rational behavior**, in particular **rational expectations**;
- In a **nonlinear** world, simple **heuristics** that work reasonably well may be the best **boundedly rational** agents can achieve